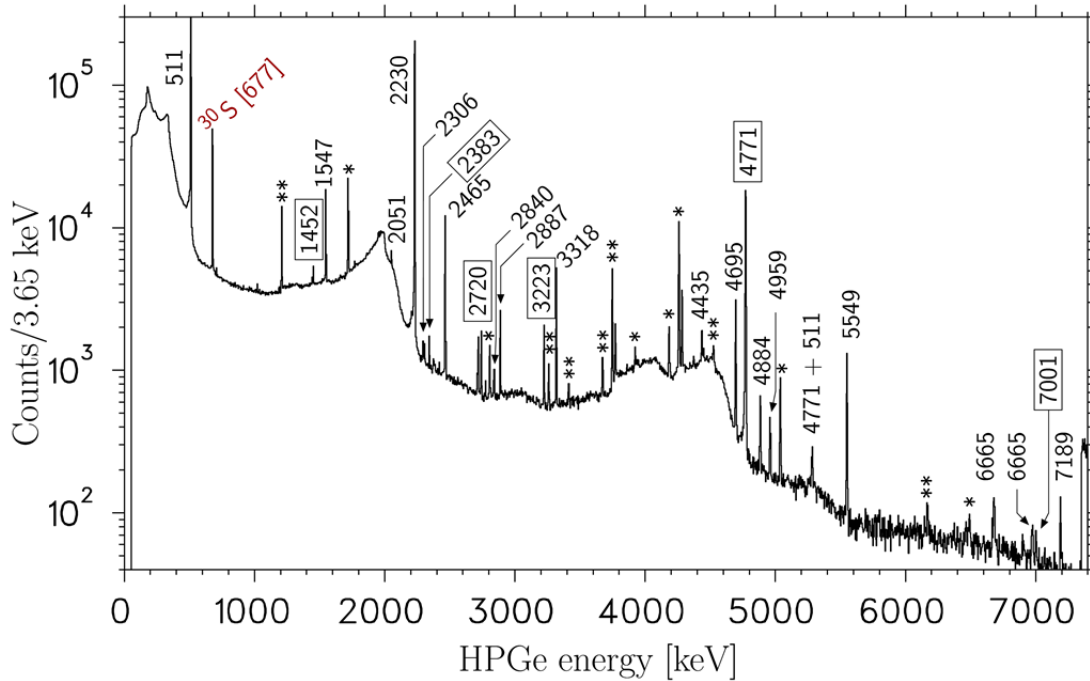


## The absolute $\gamma$ and $\beta$ branches following the $\beta^+$ decay of $^{32}\text{Cl}$

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The  $\beta$  and  $\gamma$  branches of the  $\beta^+$  decay of  $1^+$ ,  $T = 1$   $^{32}\text{Cl}$  have been measured to provide a discriminating test of the isospin-mixing corrections,  $\delta_C$ , applied to the  $0^+$  to  $0^+$  superallowed transitions used to determine  $V_{ud}$ . Up until now, there have been no nuclei studied where  $\delta_C$  is larger than  $\sim 2\%$ . We focused on improving the branching ratio of the transition to the analogue  $1^+$ ,  $T = 1$  state in  $^{32}\text{S}$  whose position in the spectrum at 7002-keV excitation is very close to a known  $1^+$ ,  $T = 0$  state at 7190 keV; this greatly enhances the size of the isospin-symmetry-breaking (ISB) correction. Our calculation of the ISB effect for this nucleus is  $\delta_C = 4.6(5)\%$  [1], a value significantly larger than those found in any of the nuclei used to extract  $V_{ud}$ . In this way,  $^{32}\text{Cl}$  provides a unique opportunity to test ISB calculations where the effect is relatively quite large. Furthermore, a precise knowledge of the absolute  $\beta$  and  $\gamma$  branches following the  $\beta$  decay of  $^{32}\text{Cl}$  may be used as an *in situ* efficiency calibration for the  $\beta$ -delayed proton decay of  $^{32}\text{Ar}$  which itself may be used to test isospin-mixing corrections (see Ref. [2]).

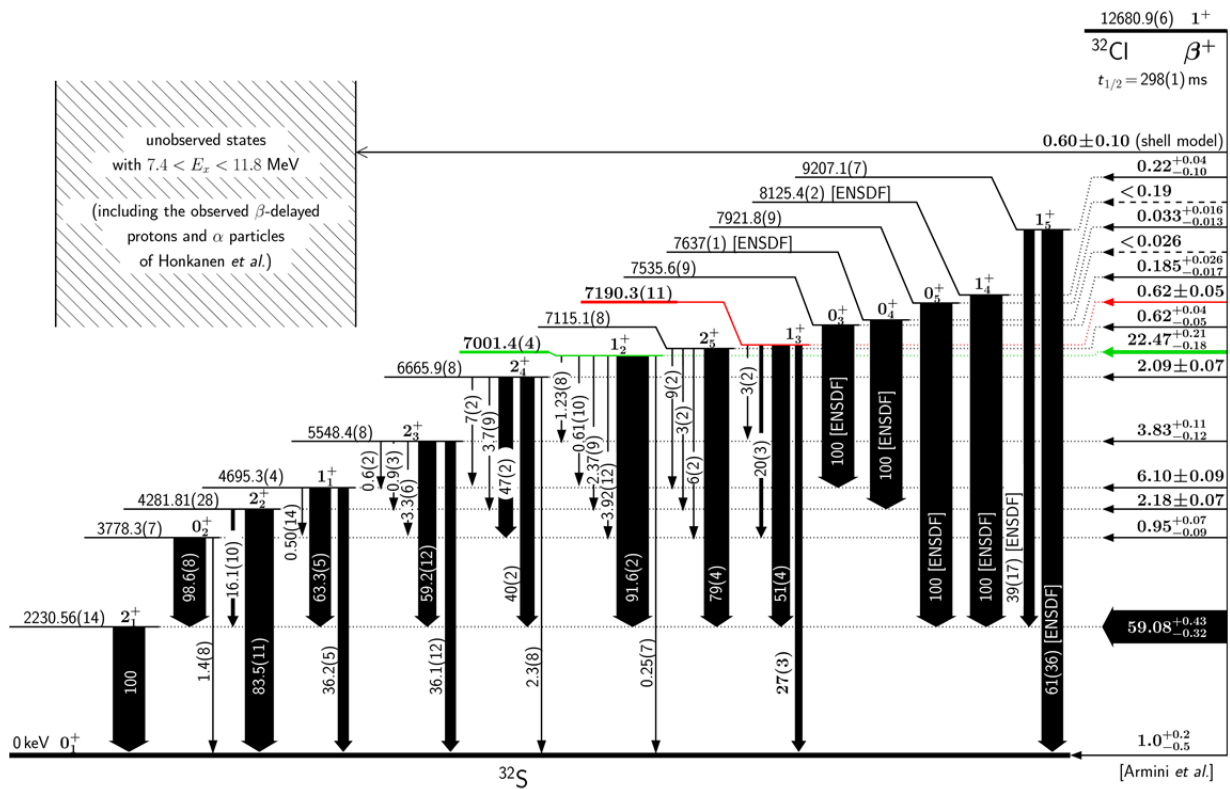
The  $^{32}\text{Cl}$  was produced via the inverse-kinematic transfer reaction  $^1\text{H}(^{32}\text{S}, n)^{32}\text{Cl}$  using a LN<sub>2</sub>-cooled, H<sub>2</sub> gas target with a 400 nA  $^{32}\text{S}$  primary beam at 24.8 MeV/nucleon. The reaction products were spatially separated by MARS, resulting in a 91% pure, 20 MeV/nucleon  $^{32}\text{Cl}$  beam with an intensity of  $\sim 2 \times 10^5$  ions/s. The beam was implanted and collected in an aluminized-Mylar tape for 0.8 s before a fast tape-transport system moved the activity to a shielded counting station 90 cm away. Data for  $\beta$ - $\gamma$  coincident events, shown in Fig. 1, were acquired using a 1.5 inch-diameter, 1 mm-thick scintillator and a



**FIG. 1.**  $\gamma$  spectrum observed by the HPGe in coincidence with a  $\beta$ . Full-energy peaks are labeled with their energy, and ones associated with the 7001-keV isobaric analogue state are highlighted with boxed values. The symbol \* and the symbol \*\* refer to single- and double-escape peaks, respectively.

70% HPGe detector. Count times were for 1, 2, and 4 sec (83%, 11%, and 6% of the data, respectively). The scintillator was placed 0.5 cm from the activity, detecting >40 keV positrons with ~32% efficiency. On the opposite side of the tape the HPGe was placed a large distance away (15.1 cm) to reduce the effects of coincidence summing of the  $\gamma$  rays. The cycle of collecting, transporting and measuring the  $^{32}\text{Cl}$  activity was repeated continuously throughout the experiment.

Once the peak areas were obtained, we used the precisely known efficiency of the HPGe detector to convert them into relative yields of  $\gamma$  rays. We then fit the  $\beta$  and  $\gamma$  branches to reproduce these yields. Our measurement has found 3 new  $\beta$  branches, 22 new  $\gamma$  lines, placed limits on 10 potential  $\gamma$  transitions, and improved the precision of the branches and yields reported previously [3, 4] by about an order of magnitude and is summarized in Fig. 2.



**FIG. 2.** Decay scheme for  $^{32}\text{Cl}$  based on this work (with results taken from Armini et al. and ENSDF labeled). The  $1^+$ ,  $T = 1$  analogue state at 7001 keV mixes with the nearby  $1^+$ ,  $T = 0$  state, giving rise to the (relatively) large observed isospin-mixing correction.

With regard to ISB effects, we found the  $\beta$  branch to the isobaric-analogue state (IAS) at 7002 keV is  $R = (22.47 \pm 0.13^{+0.16}_{-0.12})\%$ . The first uncertainty is statistical and the second is dominated by two sources of systematic uncertainty:  $^{+0.11}_{-0.12}\%$  is from the  $(1.0^{+0.2}_{-0.5})\%$  ground state branch reported by Armini et al. [5], and  $\pm 0.10\%$  from the photopeak efficiency of the HPGe detector. The partial half-life of the superallowed branch is given by  $t = \frac{t_{1/2}}{R}(1 + P_{\text{EC}})$ , where the  $^{32}\text{Cl}$  half-life is  $t_{1/2} = 298(1)$  ms,  $R$  is the superallowed branching ratio quoted above, and the small electron-capture fraction is calculated to be  $P_{\text{EC}} = 0.071\%$ . The statistical rate function is calculated to be  $f = 2411.6 \pm 2.3 \pm 0.3$ , where the first uncertainty

is from the  $Q_{EC}$  value and the second is from the shell-model calculation of the shape-correction factor. Combined, we find the experimental  $ft$  value for decay to the IAS is  $ft = 3200(30)$  s, where the precision is dominated by the  $\pm 0.9\%$  uncertainty in the branch to the IAS. This result can be interpreted in terms of an ISB effect via

$$\delta_C^{exp} = 1 + \delta_{NS} - \frac{2\langle Ft^{0^+ \rightarrow 0^+} \rangle}{ft(1 + \delta'_R)[B(F) + B(GT)]},$$

where  $\delta_{NS}$  is a nuclear-structure-dependent radiative correction calculated to be  $-0.15(2)\%$ ,  $\langle Ft^{0^+ \rightarrow 0^+} \rangle = 3071.81(83)$ s is the average corrected  $ft$  value of the 13 most precisely measured pure Fermi superallowed transitions [6],  $\delta'_R = 1.421(32)\%$  is another radiative correction, and  $B(GT)$  is calculated to be  $(1.8_{-1.7}^{+2.3}) \times 10^{-3}$  which is negligibly small compared to the dominant Fermi strength of  $B(F) = 2$ . Thus we arrive at an ISB effect of  $\delta_C^{exp} = 5.3(9)\%$ , the largest ever determined in a superallowed Fermi transition.

This result has been recently published in Refs. [7, 8].

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